AST SPECIAL TOPICS: EXOPLANETS MISC. LECTURE NOTES: SPRING 2018

INSTRUCTORS: JOEL KASTNER & MICHAEL RICHMOND

WEEKS 10 AND 11: GIANT (AND TERRESTRIAL) PLANET FORMATION & PLANET-DISK INTERACTIONS

- Giant planet formation: overview of D'Angelo, Durison, & Lissauer chapter in *Exoplanets* (p. 319), supplemented by selected M. Wyatt slides ("4. planet formation")
- Again, planets form in disks (see *Wyatt slide 6*). Two "competing" models: (1) core accretion and (2) gravitational instability
 - (1) **Core accretion** begins w/ terrestrial-planet-formation-like process buildup of planetesimals from dust and is followed by accretion of gaseous envelope from protoplanetary disk.
 - Dust grains coagulate into larger particles (Wyatt slide 7), which settle to disk midplane (Wyatt slide 9)
 - Grain coagulation process may be accelerated if grains develop "mantles" (coatings) of volatile ices (H₂O, CO)...hence observers are in hot pursuit of evidence for "snow lines" in disks
 - cm-sized particles eventually (somehow!) aggregate into km-size bodies: planetesimals (Wyatt slides 10, 11)
 - >km-sized planetesimals are compacted by their own gravity; can transition from "orderly growth," sweeping up disk material along orbit, to "runaway growth" phase, involving gravitational focusing (Wyatt slide 17)
 - Planetesimals grow into embryos via pair-wise collisions (Wyatt slide 20); larger embryos — Wyatt uses accepted term "oligarchs," but I prefer Big Mamma planetesimals — tend to sweep up all smaller planetesimals in their orbital region. Terrestrial planet formation: no need to read further; stop here.
 - When escape velocity from surface of embryo exceeds local thermal speed of disk gas, the gas can accrete onto the embryo — we would then call

INSTRUCTORS: JOEL KASTNER & MICHAEL RICHMOND

this embryo a giant planet *core*, and the accreted gas begins to form an atmosphere, and eventually, its *envelope* (*Wyatt slide 24*)

- If protoplanet's radiation trapping becomes efficient, then it can't inhibit further accretion; pressure no longer balances gravitational force, and the envelope contracts rapidly \rightarrow envelope "collapse;" happens when $M_c \approx M_e$ (Wyatt slide 24)
- Above "feedback loop" facilitates rapid accretion; planet is now in "runaway accretion" phase, regulated only by available disk gas in its vicinity. Speculation (beware): perhaps for Jupiter, Saturn lots of disk gas left in their vicinities after envelope collapse, resulting in gas giants for which $M_c \ll M_e$; whereas for Uranus, Neptune, very little disk gas left in their vicinities after envelope collapse, resulting in "ice giants" for which $M_c \sim M_e$.
- Even if they open a large gap in disk as a consequence of runaway accretion, giant planets can migrate, so can continue to slurp up additional disk gas (*Wyatt slide 26*).
- Kley & Nelson's ARAA review, "Planet-Disk Interaction and Orbital Evolution" (Kley & Nelson 2012, ARAA, 50, 211),
 http://www.annualreviews.org/doi/pdf/10.1146/annurev-astro-081811-125523 is an excellent, very dense review of planet migration theory & simulations. We will just discuss/ponder Fig. 1 and briefly review the three main types of planet migration in disks:
 - **Type I migration:** Applies to low-mass planet in a relatively massive disk. Planet's mass is too small to affect disk dynamically; in particular, the planet cannot open a disk gap. Migration is then determined by the competition between (a) *Lindblad torques* caused by resonance-driven disk spiral density waves and (b) *corotation torques* caused by disk material that is (nearly) corotating with the planet. Bottom line: complicated & messy math! Can result in either inward or outward migration.
 - **Type II migration:** Applies to massive planets that are capable of opening gaps in disks. A gap forms because, quoting Kley & Nelson, "material inside (outside) the planet loses (gains) angular momentum and recedes from the planet. Consequently, the material appears to be pushed away from the location of the planet and a gap begins to open in the disk." Lindblad and corotation torques now no longer matter. Instead, "the planet is coupled to the viscous evolution of the disk." My translation: the disk slowly accretes onto the central star, and the planet goes along for the ride!

- **Type III migration:** "Runaway" migration whose rate increases due to the effects of the migration itself on the disk. This sounds like a feedback loop, which it is. Can be inward or outward.
- See §2.4.3 of Kley & Nelson for a planet formation/migration scenario that invokes all three of the above migration types...but not in I–II-III order!
- (2) Gravitational instability (GI) models of giant planet formation in dusty molecular disks were developed via analogy to star formation in dusty molecular clouds: gas-phase fragmentation of the disk into bound clumps (Boss 1997).
 - GIs build out of local perturbations in steady-state disk conditions (density, temperature) (Wyatt slides 41-43)
 - Stability to perturbations parameterized through Toomre Q:

$$Q = \frac{c\kappa}{\pi G \Sigma}$$

where $c = (\gamma kT/\mu m_H)^{1/2}$ is the local sound speed, κ is oscillation frequency of a test particle or parcel of gas about its equilibrium position — for a disk, $\kappa = \Omega$, i.e., the Keplerian angular velocity — and Σ is local surface density.

If Q < 1 then the disk is locally unstable to collapse.

- Note that this condition implies one wants a cool disk with a high surface density. So conditions are most favorable for small Q in massive disks; conditions for small Q are even favorable in outer regions of massive disks. So...
- GI models predict planet formation (a) should occur rapidly and (b) may occur at large radii.
- Both of the foregoing predictions are supported by ALMA images of very young disks, like HL Tau (whose age is likely < 1 Myr):
 http://www.almaobservatory.org/en/press-release/revolutionary-alma-image-reveal *Caveat:* There are other mechanisms that could explain gaps in disks: snow lines, density waves, magnetic fields(!)...

TABLE 1. Comparison: giant planet formation models

	Core Accretion	Gravitational Instability
Timescale	Myr $(10^4 - 10^5 \text{ orbital periods})$	kyr (tens of orbital periods)
Disk masses	MMSN $(M_d \sim 0.01 M_{\odot})$ enough?	massive $(M_d \stackrel{>}{\sim} 0.1 M_{\odot})$
planet formation regions	a few AU to tens of AU	can extend to hundreds of AU